

# Motivating Matrix-Free Finite Elements

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Discretization of PDE for computation presents several choices

Matrix-free finite elements present benefits and challenges

We'll see how matrix-free finite elements works and the downstream choices and challenges that this discretization provides

# Overview

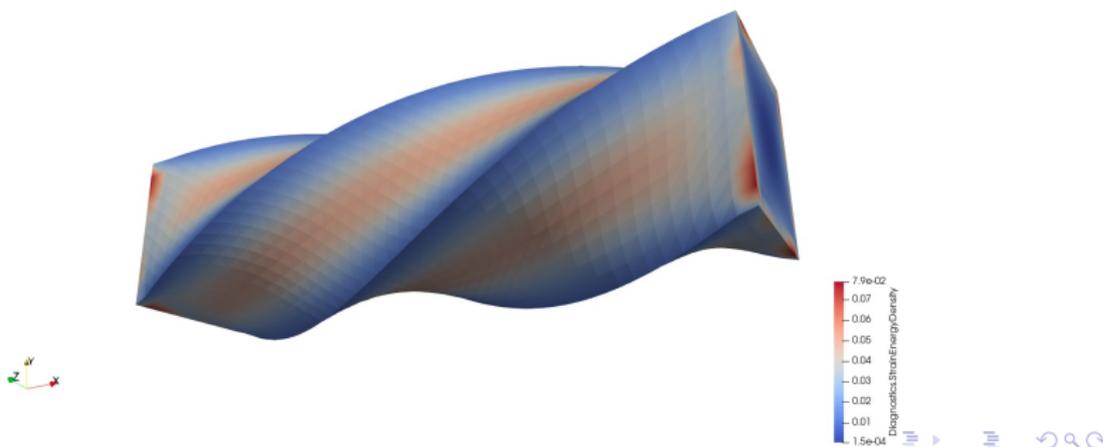
- 1 Finite Elements
- 2 Matrix-Free
- 3 Non-Linear Problem
- 4 Preconditioning
- 5 Cool Results
- 6 Questions

## End Goal

We need to transform equations:

$$\begin{aligned} -\nabla \cdot \sigma - \rho g &= 0, \text{ in } \Omega \\ u &= \bar{u}, \text{ on } \partial\Omega^D \\ \sigma \cdot n &= \bar{t}, \text{ on } \partial\Omega^N \end{aligned}$$

into simulations:



# Equation to Solve

Consider the general non-linear residual

$$F(u) - R = 0$$

where

- $F(u)$  represents our 'physics'
- $R$  is some 'right hand side' or forcing term
- $u$  is the unknown solution

# Problem

Computers like discrete equations/operations

Our equation is continuous, and computers *hate* that

# Another Problem

Solving these non-linear PDEs can be hard

Really, really hard

# First Decision

With the  $L^2$  inner product

$$\langle v, u \rangle = \int_{\Omega} vu \, dV$$

We can reformulate to the 'weak form':

$$\langle v, f(u) \rangle = \int_{\Omega} v \cdot f_0(u, \nabla u) + \nabla v : f_1(u, \nabla u) \, dV = 0$$

Note:

- boundary terms are omitted here (integration by parts)
- $v$  is in some homogeneous space  $V_0$

$f_0$  and  $f_1$  Examples

Projection problem:

$$u = g$$

$$\int_{\Omega} v \cdot u \, dV = \int_{\Omega} v g \, dV$$

$$f_0(u, \nabla u) = u, \quad f_1(u, \nabla u) = 0$$

Poisson problem:

$$-\nabla^2 u = g$$

$$\int_{\Omega} \nabla v : \nabla u \, dV - \int_{\partial\Omega} v \nabla u \cdot n \, dS = \int_{\Omega} v g \, dV$$

$$f_0(u, \nabla u) = 0, \quad f_1(u, \nabla u) = \nabla u$$

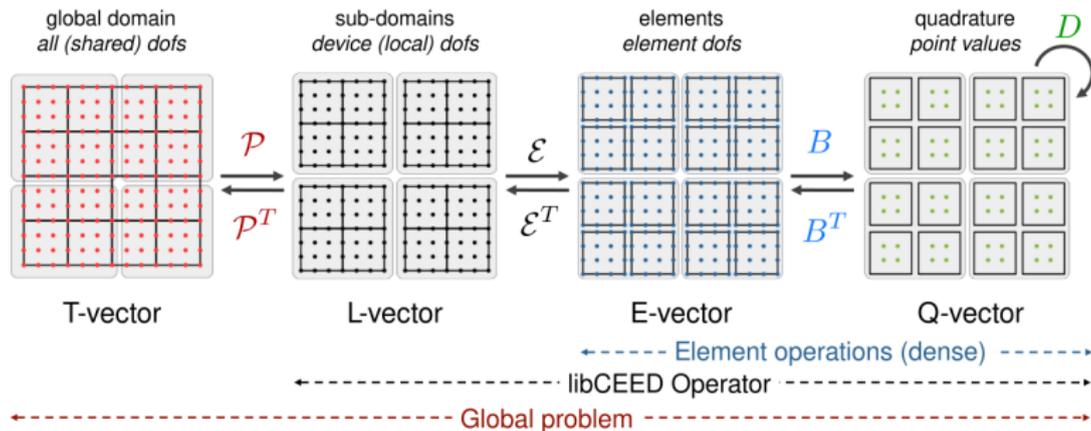
## Second Decision

How to discretize?

- $u, v \in V$
- Different  $V$  give different properties
- We select  $V$  to be the space of piecewise continuous polynomials
- No continuity in derivatives (not splines)

## Second Decision

$$A = \mathcal{P}^T \mathcal{E}^T B^T D B \mathcal{E} \mathcal{P}$$



T(otal)-vector represents discretization

Want solution values on red dots (Degrees of Freedom [DoF])

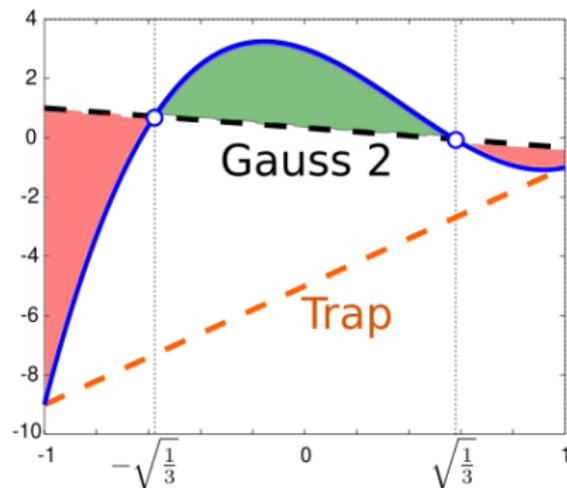
Fitting piecewise polynomial to red dots gives continuous solution

# Third Decision

Need to convert integration to a discrete problem

- Quadrature rules take values at points to compute integral
- Piecewise nature means each element can be integrated separately
- Workflow:
  - 1 Map to interpolating polynomial
  - 2 Evaluate values/derivatives at quadrature points
  - 3 Evaluate  $f_0, f_1$
  - 4 Multiply by quadrature weights
  - 5 Transpose process to return to DoFs

## Third Decision



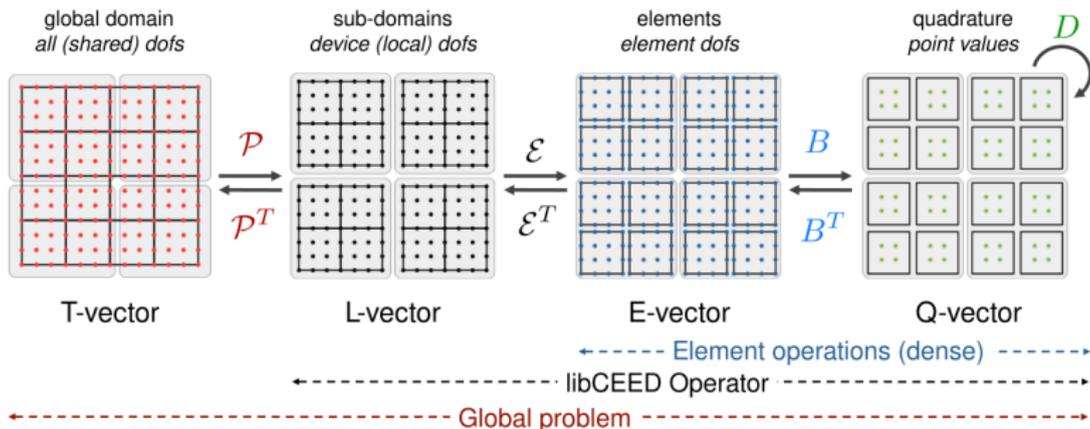
Why the extra work of interpolating to quadrature points?

Better accuracy in discrete integration

Gauss-Legendre accurate for polynomials up to  $2n - 1$

## Third Decision

$$A = \mathcal{P}^T \mathcal{E}^T B^T D B \mathcal{E} \mathcal{P}$$



E(lement)-vector represents each 'piece' of our solution polynomial

DoFs on blue dots are red dots on a specific element

$B$  maps to the quadrature points - green dots

$D$  evaluates  $f_0, f_1$ , multiply by quadrature weights

# Galerkin Form

Galerkin form for an arbitrary second order PDE:

$$\sum_e \mathcal{E}^T \left[ (\mathbf{N}^e)^T \mathbf{W}^e \Lambda (f_0(u^e, \nabla u^e)) + \sum_{i=0}^{d-1} (\mathbf{D}_i^e)^T \mathbf{W}^e \Lambda (f_1(u^e, \nabla u^e)) \right] = 0$$

- $\mathcal{E}$  - element assembly/restriction operator
- $\mathbf{N}^e$  - interpolation to quadrature points
- $\mathbf{D}_i^e$  - derivatives at quadrature points
- $\mathbf{W}^e$  - quadrature weights
- $\Lambda$  - pointwise multiplication at quadrature points
- $u^e = \mathbf{N}^e \mathcal{E}^e u$  and  $\nabla u^e = \{\mathbf{D}_i^e \mathcal{E}^e u\}_{i=0}^{d-1}$

## Wait, What?!?

Non-linear problem:

$$F(u) - R = 0$$

Weak form:

$$\langle v, f(u) \rangle = \int_{\Omega} v \cdot f_0(u, \nabla u) + \nabla v : f_1(u, \nabla u) dV = 0$$

Galerkin form:

$$\sum_e \mathcal{E}^T \left[ (\mathbf{N}^e)^T \mathbf{W}^e \Lambda(f_0(u^e, \nabla u^e)) + \sum_{i=0}^{d-1} (\mathbf{D}_i^e)^T \mathbf{W}^e \Lambda(f_1(u^e, \nabla u^e)) \right] = 0$$

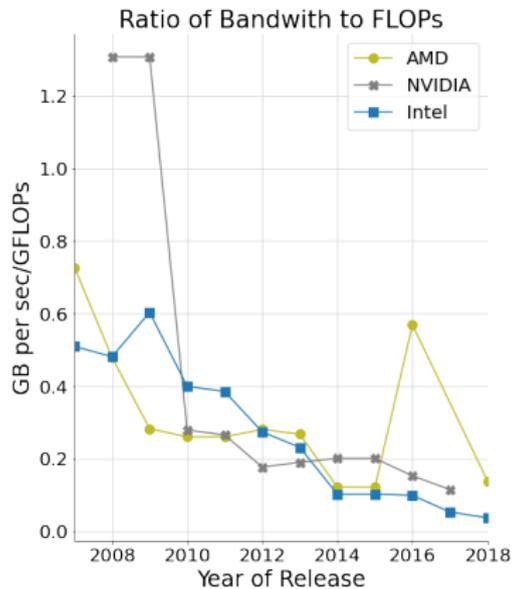
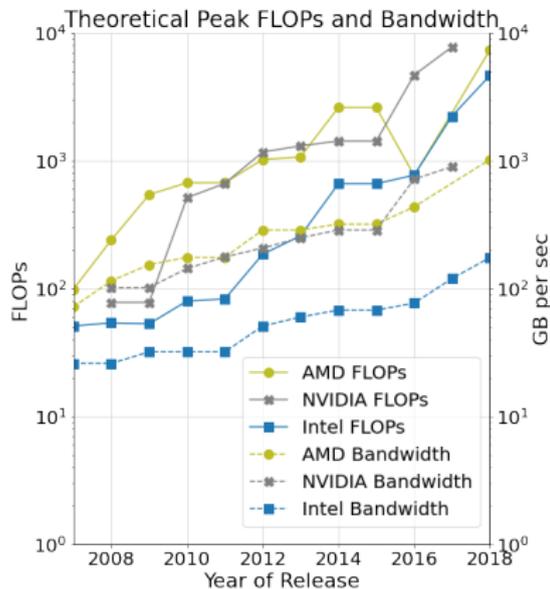
# What I'm Handwaving

- Elements can have deformed shapes (unstructured meshes)
- Need 'geometric' factors for each individual element  
(Change of coordinate systems from physical to reference)
- $\mathcal{P}$  is parallel communication (supercomputers!)
- $R$  (and its weak analogue) can be considered part of  $f_0$  or separate
- Boundary terms are similar but 1 dimension lower (i.e.  $3D \rightarrow 2D$ )

# Traditional Approach

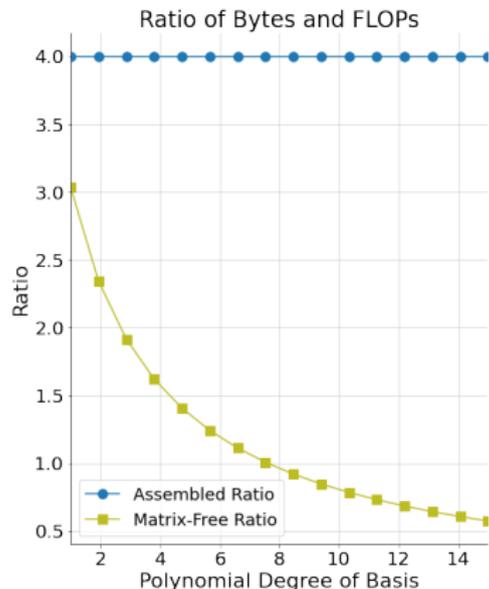
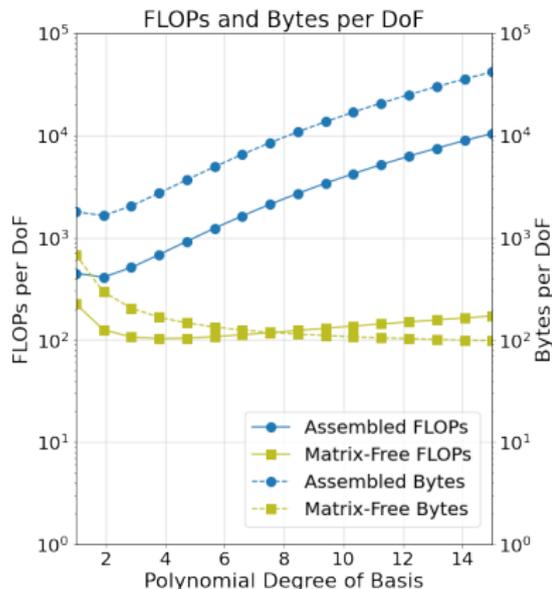
- The Galerkin form naturally creates a sparse matrix for linear  $F(u)$
- This sparse matrix can be solved with linear algebra
- Two problems:
  - 1  $F(u)$  is not usually linear
  - 2 Sparse matrices don't use hardware optimally
- Big Idea: perform the operations without forming a matrix

# Modern Hardware



Modern hardware has lower memory bandwidth than FLOPs

## Matrix-Free



Requirements for matrix-vector product with sparse matrix vs matrix-free  
for screened Poisson  $\nabla^2 u - \alpha^2 u = f$  in 3D

**Matrix-free representations using tensor product bases  
better match modern hardware limitations**

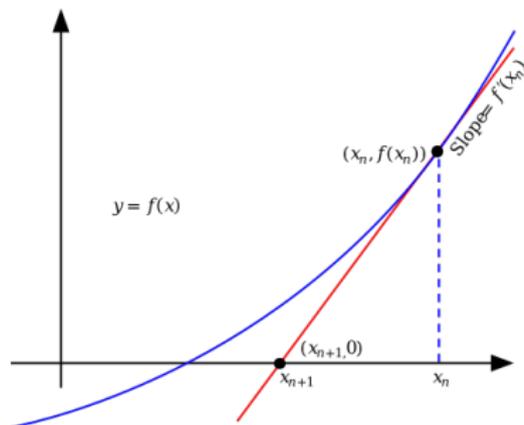
# Non-linear Problem

We still have a non-linear problem

Linear algebra solves *linear* problems

# Newton's Method

Many non-linear solvers fit into the 'Newton' family



In Newton's method, we form a tangent line

The *linear* equation is easy to solve

In higher dimensions, the line becomes a (hyper)plane

# Newton's Method

We want to solve

$$F(u) - R = 0$$

We can instead solve a series of 'tangent plane' problems

$$F(u_k) - R + J(u_k)(u_k - u_{k+1}) = 0$$

$u_{k+1}$  provides closer estimates of the true solution  $u_*$

We'll denote  $w = u_k - u_{k+1}$

# A New Problem

Ok, but we still need to form the tangent plane  $J(u_k)$

This feels like we keep changing what the problem is!

## Beauty

Weak form:

$$\langle v, f(u) \rangle = \int_{\Omega} v \cdot f_0(u, \nabla u) + \nabla v : f_1(u, \nabla u) dV = 0$$

We want to take the derivative of this weak form

Note:

- $\frac{d}{dx} (f(x) + g(x)) = f'(x) + g'(x)$
- $\frac{d}{dx} f(g(x)) = f'(g(x)) g'(x)$

## Beauty

Weak form derivative:

$$\langle v, J(u_k) w \rangle = \int_{\Omega} \begin{bmatrix} v^T & (\nabla v)^T \end{bmatrix} \begin{bmatrix} f_{0,0} & f_{0,1} \\ f_{1,0} & f_{1,1} \end{bmatrix} \begin{bmatrix} w \\ \nabla w \end{bmatrix}$$

Note:

- $f_{i,0} = \frac{\partial f_i}{\partial u}(u_k, \nabla u_k)$
- $f_{i,1} = \frac{\partial f_i}{\partial \nabla u}(u_k, \nabla u_k)$
- Can be discretized the same way as before

# Wait, What?!? - Non-linear Equation

Non-linear problem:

$$F(u) - R = 0$$

Weak form:

$$\langle v, f(u) \rangle = \int_{\Omega} v \cdot f_0(u, \nabla u) + \nabla v : f_1(u, \nabla u) dV = 0$$

Galerkin form:

$$\sum_e \mathcal{E}^T \left[ (\mathbf{N}^e)^T \mathbf{W}^e \Lambda(f_0(u^e, \nabla u^e)) + \sum_{i=0}^{d-1} (\mathbf{D}_i^e)^T \mathbf{W}^e \Lambda(f_1(u^e, \nabla u^e)) \right] = 0$$

# Wait, What?!? - Linear Equation

Linearized problem:

$$F(u_k) - R + J(u_k)(u_k - u_{k+1}) = 0$$

Weak form:

$$\langle v, J(u_k) w \rangle = \int_{\Omega} \begin{bmatrix} v^T & (\nabla v)^T \end{bmatrix} \begin{bmatrix} f_{0,0} & f_{0,1} \\ f_{1,0} & f_{1,1} \end{bmatrix} \begin{bmatrix} w \\ \nabla w \end{bmatrix}$$

Galerkin form:

$$\sum_e \mathcal{E}^T \left[ (\mathbf{N}^e)^T \mathbf{W}^e \Lambda(f_0^*(w^e, \nabla w^e)) + \sum_{i=0}^{d-1} (\mathbf{D}_i^e)^T \mathbf{W}^e \Lambda(f_1^*(w^e, \nabla w^e)) \right] = 0$$

# A New Problem

We have a discrete system we can solve efficiently

But we still have a problem

(Our problems are getting easier to tackle though!)

# Condition Number

Matrix-free representation of the linear problem restricts solvers

Krylov subspace methods only need the matrix-vector product  $Ax$ ,  
not the assembled matrix  $A$

But Krylov subspace methods are sensitive to condition number

High-order finite elements are more efficient matrix-free,  
but high-order finite elements have a higher condition number

# Preconditioning

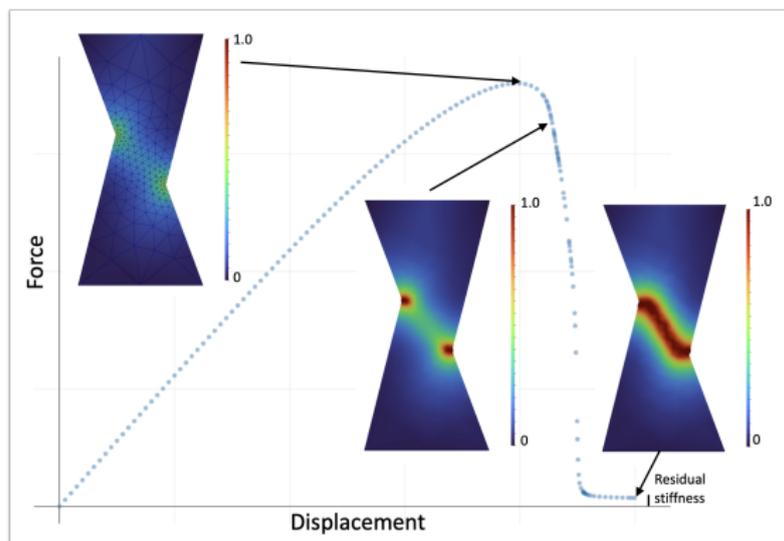
Preconditioning improves the condition number of  $A$

Big Idea:  $P^{-1}A \approx I$  has a better condition number

Challenge: Finding efficient  $P^{-1}$  can be difficult

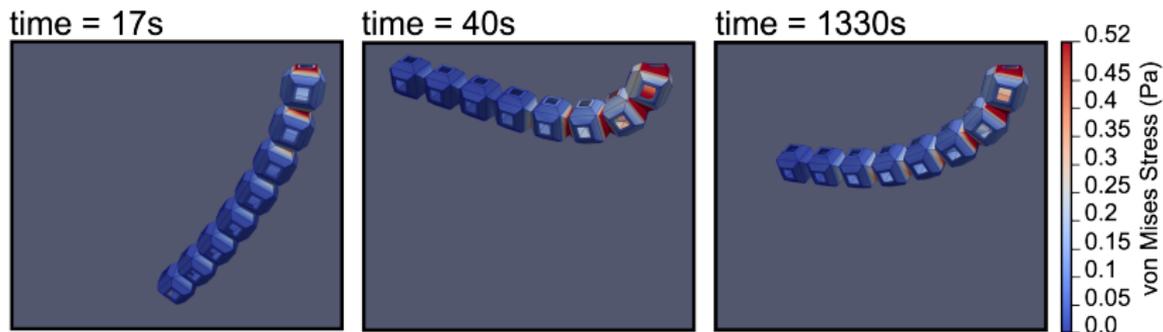
The Fix: This is a rich research area with lots of options

# Example - Brittle Fracture



Quasistatic simulation of a compressive shear test for a generic brittle material

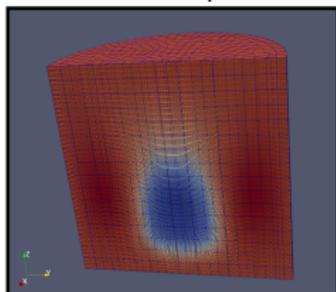
# Example - Schwarz Pendulum



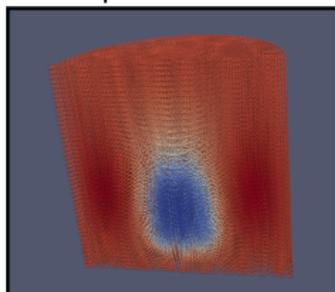
Dynamic simulation of 8 Schwarz-P cells with an applied traction force creating pendulum-like motion

# Example - Sinker Simulation

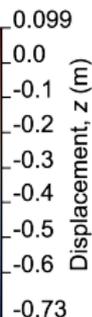
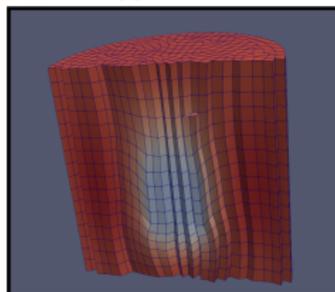
iMPM - mesh &amp; particles



iMPM - particles

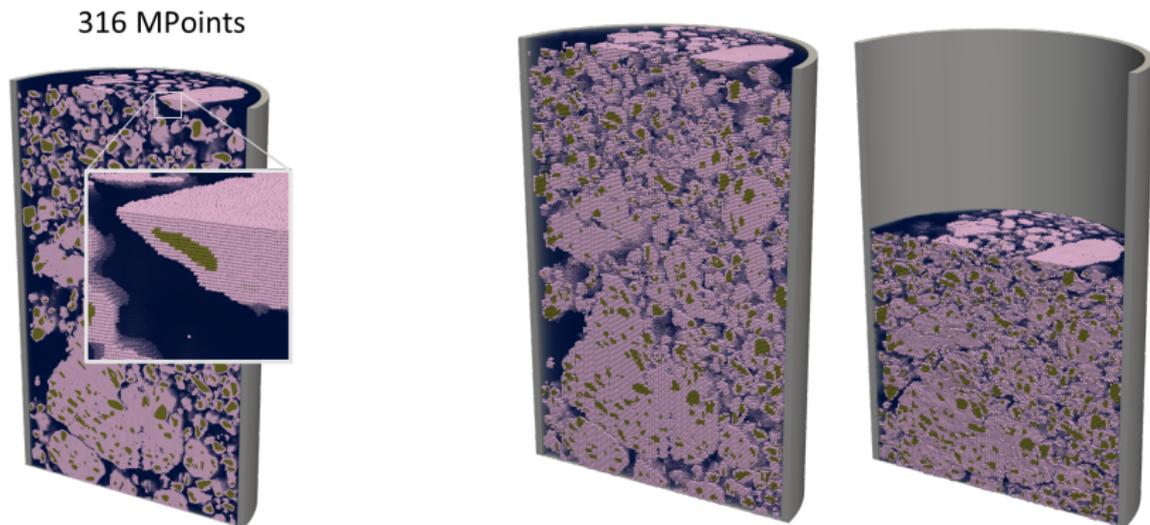


FEM - mesh



Quasistatic simulation of a cylinder with a dense high-modulus inclusion surrounded by a low-density low-modulus near-incompressible "foam"

# Example - Press Simulation



Compression of mock HE grains (gold) and binder (pink) mixture  
(Reset background mesh to computational region on each timestep)

# Questions?



libCEED Repo: <https://github.com/CEED/libCEED>

Ratel Repo: <https://gitlab.com/micromorph/ratel>

Grant: Predictive Science Academic Alliance Program (DE-NA0003962)



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